

# FRACTIONAL POISSON FIELDS

NIKOLAI LEONENKO AND ELY MERZBACH

ABSTRACT. Using inverse subordinators and Mittag-Leffler functions, we present a new definition of a fractional Poisson process parametrized by points of the Euclidean space  $\mathbb{R}_+^2$ . Some properties are given and, in particular, we prove a long-range dependence property.

## 1. INTRODUCTION

There are essentially four different approaches to the concept of Fractional Poisson process on the real line: The “integral representation” method follows the integral representation of the Fractional Brownian motion, replacing the Brownian motion by the Poisson process (Wang, Wen [32]). For some other aspects of this approach, see [14],[6].

Another approach that we can call the “Renewal” approach consists of considering the characterization of the Poisson process as a sum of independent non-negative random variables, and instead of the assumption that these random variables have an exponential distribution, we assume that they have the Mittag-Leffler distribution (Mainardi et al. [19],[20], Repin, Saichev [28]). In [3], the renewal approach to the fractional Poisson process is developed and it is proved that its one-dimensional distributions coincide with the solution to fractionalized state probabilities.

A third approach, the “differential” one, uses the differential equations of the Poisson process and replaces them by fractional derivatives (Beghin, Orsingher [3],[4]). Finally, using “inverse subordinator”, a kind of fractional Poisson process can be constructed (Meerschaert, Nane, Vellaisamy [23]). Many formulas attributed to [23] have been obtained in [14], [3],[4].

Here, we will follow the fourth method to generalize and define a Fractional Poisson field parametrized by points of the Euclidean space  $\mathbb{R}_+^2$ , as it has been done for fractional Brownian fields, see, e.g., [10], [18] (see also [11],[12]).

The starting point of our extension will be the set-indexed Poisson process which is a well-known concept, see, e.g., [29] or [24],[25].

---

2010 *Mathematics Subject Classification.* 60G22, 60G60, 60G55.

*Key words and phrases.* Poisson fields, long-range dependence, subordinator, inverse subordinator.

Nikolai Leonenko and Ely Merzbach were partially supported by a grant of the Commission of the European Communities PIRSES-GA-2008-230804 (Marie Curie) “Multi-parameter Multi-fractional Brownian Motion”.

## 2. SUBORDINATORS AND INVERSE SUBORDINATORS

This section collects some known results from the theory of subordinators and inverse subordinators, see Bingham [5], Veillette and Taqqu [30], [31], [22], among the others.

**2.1. Inverse subordinator.** Let  $L(t)$ ,  $t \geq 0$  be a Lévy subordinator with Laplace exponent

$$(1) \quad \phi(s) = \mu s + \int_{(0,\infty)} (1 - e^{-sx}) \Pi(dx), \quad s \geq 0,$$

where  $\mu \geq 0$  is the drift and the Lévy measure  $\Pi$  on  $\mathbb{R}_+ \cup \{0\}$  satisfies

$$(2) \quad \int_0^\infty (1 \wedge x) \Pi(dx) < \infty.$$

This means that

$$(3) \quad \mathbb{E}e^{-sL(t)} = e^{-t\phi(s)}, \quad s \geq 0.$$

Consider the inverse subordinator  $Y(t)$ ,  $t \geq 0$ , which is given by the first-passage time of  $L$

$$(4) \quad Y(t) = \inf \{u \geq 0 : L(u) > t\}, \quad t \geq 0.$$

The process  $Y(t)$ ,  $t \geq 0$ , is non-decreasing and its sample paths are a.s. continuous if  $L$  is strictly increasing. Also  $Y$  is, in general, non-Markovian with non-stationary and non-independent increments.

We have

$$(5) \quad \{L(u_i) < t_i, i = 1, \dots, n\} = \{Y(t_i) > u_i, i = 1, \dots, n\}.$$

Let

$$(6) \quad U(t) = \mathbb{E}Y(t)$$

be the renewal function of  $Y(t)$ , and let

$$(7) \quad H_u(t) = \mathbb{P}\{Y(u) < t\}.$$

From (1) we have

$$(8) \quad \int_0^\infty e^{-st} dH_u(t) = e^{-t\phi(s)},$$

and

$$(9) \quad \tilde{U}(s) = \int_0^\infty U(t) e^{-st} dt = \frac{1}{s\phi(s)};$$

thus,  $\tilde{U}$  characterizes the process  $Y$  (since  $\phi$  characterizes  $L$ ).

Therefore, we get a covariance formula ([30], [31])

$$(10) \quad \text{Cov}(Y(t_1), Y(t_2)) = \int_0^{t_1 \wedge t_2} (U(t_1 - \tau) + U(t_2 - \tau)) dU(\tau) - U(t_1)U(t_2).$$

The most important example is considered in the next section, but there are some others.

**2.2. Inverse stable subordinator.** Let  $L_\alpha(t), t \geq 0$ , be a  $\alpha$ -stable subordinator with  $\phi(s) = s^\alpha, 0 < \alpha < 1$  (cadlag, continuous in probability, with independent and stationary increments), whose density  $g(t, x)$  is such that  $L_\alpha(1)$  has pdf

$$(11) \quad g_\alpha(x) = g(1, x) = \frac{1}{\pi} \sum_{k=1}^{\infty} (-1)^{k+1} \frac{\Gamma(\alpha k + 1)}{k!} \frac{1}{x^{\alpha k + 1}} \sin(\pi k \alpha) = \frac{1}{x} W_{-\alpha, 0}(-x^{-\alpha}), \quad x > 0,$$

and with Laplace transform

$$E e^{-s L_\alpha(t)} = \exp\{-t s^\alpha\}, \quad s \geq 0.$$

Here we use the Wright's generalized Bessel function (see, e.g., [9])

$$W_{\gamma, \beta}(z) = \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(1+k)\Gamma(\beta+\gamma k)}, \quad z \in \mathbb{C},$$

where  $\gamma > -1$ , and  $\beta \in \mathbb{R}^1$ , and for  $\beta = 0, \gamma = -\alpha \in (-1, 0)$

$$(12) \quad W_{-\alpha, 0}(z) = \sum_{k=1}^{\infty} \frac{\sin(\pi k \alpha)}{\pi} \frac{z^k \Gamma(1 + \alpha k)}{k!}$$

by reciprocity relation for the  $\Gamma$ -function.

Also

$$g_\alpha(x) \sim \frac{\left(\frac{\alpha}{x}\right)^{\frac{2-\alpha}{2(2-\alpha)}}}{\sqrt{2\pi\alpha(1-\alpha)}} \exp\left\{-\left(1-\alpha\right)\left(\frac{x}{\alpha}\right)^{-\frac{\alpha}{1-\alpha}}\right\}, \quad x \rightarrow 0,$$

and

$$g_\alpha(x) \sim \frac{\alpha}{\Gamma(1-\alpha)x^{1+\alpha}}, \quad x \rightarrow \infty.$$

Then

$$Y_\alpha(t) = \inf\{u \geq 0 : L_\alpha(u) > t\}$$

has density

$$(13) \quad f_\alpha(t, x) = \frac{t}{\alpha} x^{-1-\frac{1}{\alpha}} g_\alpha(tx^{-\frac{1}{\alpha}}), \quad x \geq 0, \quad t > 0.$$

Let

$$E_\alpha(z) = \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(\alpha k + 1)}, \quad \alpha > 0, \quad z \in \mathbb{C}$$

be the Mittag-Leffler function [8],[9], and recall the following:

i) The Laplace transform of the Mittag-Leffler function is of the form

$$\int_0^\infty e^{-st} E_\alpha(-t^\alpha) dt = \frac{s^{\alpha-1}}{1+s^\alpha}, \quad 0 < \alpha < 1, \quad t \geq 0.$$

(ii) The Mittag-Leffler function is a solution of the fractional equation with fractional Caputo-Djrbashian derivative  $D_t^\alpha$

$$D_t^\alpha E_\alpha(at^\alpha) = aE_\alpha(at^\alpha).$$

**Proposition.**

i)

$$Ee^{-sY_\alpha(t)} = \sum_{n=0}^{\infty} \frac{(-st^\alpha)^n}{\Gamma(\alpha n + 1)} = E_\alpha(-st^\alpha),$$

(ii) Both processes  $L_\alpha(t), t \geq 0$  and  $Y_\alpha(t)$  are self-similar

$$\frac{L_\alpha(at)}{a^{1/\alpha}} \stackrel{d}{=} L_\alpha(t), \quad \frac{Y_\alpha(at)}{a^\alpha} \stackrel{d}{=} Y_\alpha(t), \quad a > 0.$$

(iii)

$$\frac{\partial E(Y_\alpha(t_1) \cdots Y_\alpha(t_k))}{\partial t_1 \cdots \partial t_k} = \frac{1}{\Gamma^k(\alpha)} \frac{1}{[t_1(t_2 - t_1) \cdots (t_k - t_{k-1})]^{1-\alpha}}, \quad 0 < t_1 < \cdots < t_k.$$

In particular,

(A)

$$EY_\alpha(t) = \frac{t^\alpha}{\Gamma(1+\alpha)};$$

(B)

$$\text{Cov}(Y_\alpha(t), Y_\alpha(s)) =$$

$$(14) \quad = \frac{1}{\Gamma(1+\alpha)\Gamma(\alpha)} \int_0^{\min(t,s)} ((t-\tau)^\alpha + (s-\tau)^\alpha) \tau^{\alpha-1} d\tau - \frac{(st)^\alpha}{\Gamma^2(1+\alpha)}.$$

*Comments.* 1. Notice that this last property can be interpreted as long-range dependence.

2. There is a (complicated) form of all finite-dimensional distributions of  $Y_\alpha(t), t \geq 0$ , in the form of Laplace transforms, see [5].

**2.3. Mixture of inverse subordinators.** Different kinds of inverse subordinators can be considered.

Here we deal with a mixture of  $\alpha$ -stable subordinators. Following (1)

$$\phi(s) = \int_0^1 q(w) s^\alpha dw = \int_0^\infty (1 - e^{-sx}) l_q(x) dx,$$

where  $q(w)$  is a probability density on  $(0, 1)$ , and the density  $l_q(x)$  of the Lévy measure is given by

$$l_q(x) = \int_0^1 \frac{x^{-\alpha-1}}{\Gamma(-\alpha)} q(\alpha) d\alpha.$$

The  $\alpha$ -stable subordinator corresponds to the choice

$$q(w) = \delta(w - \alpha);$$

thus we can consider

$$q(w) = C_1 \delta(w - \alpha_1) + C_2 \delta(w - \alpha_2), C_1 + C_2 = 1,$$

with  $\alpha_1 > \alpha_2$ . This is a sum of two independent subordinators with

$$\phi(s) = C_1 s^{\alpha_1} + C_2 s^{\alpha_2}, C_1 + C_2 = k.$$

It was proved in [17] that

$$\tilde{U}(t) = \frac{1}{C_1 s^{\alpha_1} + C_2 s^{\alpha_2}}, U(t) = \frac{1}{C_1} t^{\alpha_1-1} E_{\alpha_1-\alpha_2, \alpha_1}(-C_2 t^{\alpha_1-\alpha_2}/C_1),$$

where  $E_{\alpha, \beta}(z)$  is the two-parametrical Mittag-Leffler function ([8],[9])

$$E_{\alpha, \beta}(z) = \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(\alpha k + \beta)}, \quad 0 < \alpha < 2, \quad \beta > 0, \quad z \in \mathbb{C}.$$

Another possibility is to consider the uniform mixture with  $q(x) = 1, x \in (0, 1)$ .

One can also consider the tempered stable inverse subordinator, an inverse subordinator to the Poisson process, compound Poisson process, Gamma and inverse Gaussian Lévy processes, for additional details see Veillette and Taqqu [30], [31].

### 3. HOMOGENEOUS POISSON FIELD

Let  $\mathcal{B}$  be the class of Borel subsets of  $\mathbb{R}^d$ .

It is well known (see, e.g., Stoyan, Kendall, Mecke [29]) that the Poisson field, also called the spatial Poisson process, is an integer-valued random measure  $N(A)$ ,  $A \in \mathcal{B}$ , which has the following properties:

**(P1)** for any finite collection  $\{A_1, \dots, A_n\}$  of disjoint Borel subsets

$$N(A_1), \dots, N(A_n)$$

are independent;

**(P2)** the probability distribution function of  $N(A)$  depends on the set  $A$  only through its size (or Lebesgue measure)  $|A|$ ;

**(P3)** there exists  $\lambda > 0$  such that

$$\mathbb{P}(N(A) \geq 1) = \lambda |A| + o(|A|), |A| \rightarrow 0;$$

**(P4)** the probability of points overlapping is zero

$$\lim_{|A| \rightarrow 0} \frac{\mathbb{P}(N(A) \geq 1)}{\mathbb{P}(N(A) = 1)} = 1.$$

If these axioms are satisfied, we have for any finite  $n = 1, 2, \dots$ , and for any disjoint bounded Borel sets  $A_1, \dots, A_n$

$$(15) \quad \begin{aligned} & \mathbb{P}(N(A_1) = k_1, \dots, N(A_n) = k_n) \\ &= \frac{\lambda^{k_1 + \dots + k_n}}{k_1! \cdots k_n!} (|A_1|)^{k_1} \cdots (|A_n|)^{k_n} \exp \left\{ - \sum_{j=1}^n \lambda |A_j| \right\}, \quad k_j = 0, 1, 2, \dots, \end{aligned}$$

while

$$\mathbb{E}N(A) = \lambda |A|, \quad \text{Cov}(N(A_1), N(A_2)) = \lambda |A_1 \cap A_2|.$$

Also the emptiness probabilities

$$\mathbb{P}(N(A) = 0) = e^{-\lambda |A|}.$$

Particular emptiness probabilities are given by the so-called contact distribution function

$$H_A(r) = 1 - \mathbb{P}(N(rA) = 0), \quad r \geq 0.$$

Here  $A$  is a compact ‘‘test set’’ with  $|A| > 0$  containing the origin  $0 \in A$ , and  $rA = \{rx, x \in A\}$  denotes the dilation of  $A$ . Obviously  $H_A(r)$  is a distribution function.

#### 4. FRACTIONAL POISSON FIELDS

For the sake of simplicity, we restrict our discussion to the two-dimensional case  $d = 2$ .

Following ideas of Meershaert, Nane and Vellaisamy [23] (see also Mainardi, Gorenflo, Scalas [19], Beghin and Orsingher [3], [4]), we define the fractional Poisson field as follows:

$$(16) \quad N_{FP}(t_1, t_2) = N(Y^{(1)}(t_1), Y^{(2)}(t_2)), \quad (t_1, t_2) \in \mathbb{R}_+^2,$$

where  $Y^{(1)}(t), t \geq 0$  and  $Y^{(2)}(t), t \geq 0$  are two independent inverse subordinators, and

$$N(t_1, t_2) = N([0, t_1] \times [0, t_2]), \quad (t_1, t_2) \in \mathbb{R}_+^2.$$

is a restriction of homogeneous Poisson field on  $\mathbb{R}_+^2$ . We also assume that the Poisson field  $N(t_1, t_2)$  is independent of the inverse subordinators  $Y^{(1)}(t)$  and  $Y^{(2)}(t)$ .

More precisely, if  $Y_{\alpha_1}^{(1)}(t), t \geq 0$  and  $Y_{\alpha_2}^{(2)}(t), t \geq 0$  are two independent inverse stable subordinators with indices  $\alpha_1 \in (0, 1)$  and  $\alpha_2 \in (0, 1)$ , which are independent of the Poisson field  $N(t_1, t_2)$ , then we denote the fractional Poisson process

$$(17) \quad N_{\alpha_1, \alpha_2}(t_1, t_2) = N(Y_{\alpha_1}^{(1)}(t_1), Y_{\alpha_2}^{(2)}(t_2)), \quad (t_1, t_2) \in \mathbb{R}_+^2.$$

For the dimension  $d = 1$ , this construction is known as the fractional Poisson process (see Meershaert, Nane and Vellaisamy [23]). It is denoted by

$$N_\alpha(t) = N(Y_\alpha(t)), \quad t \geq 0, \quad \alpha \in (0, 1),$$

where  $N(t) = N([0, t]), t \geq 0$ , is the classical homogeneous Poisson process with parameter  $\lambda > 0$ , which is independent of the inverse stable subordinator  $Y_\alpha(t), t \geq, \alpha \in (0, 1)$ .

One can compute the following expression for the one-dimensional distribution:

$$\begin{aligned} \mathbb{P}(N_\alpha(t) = k) &= \frac{(t\lambda)^k}{k!} \int_0^\infty e^{-\lambda x} x^{k-1-\frac{1}{\alpha}} g_\alpha(tx^{-\frac{1}{\alpha}}) dx \\ &= \frac{(\lambda t^\alpha)^k}{k!} \sum_{j=1}^\infty \frac{(k+j)!}{j!} \frac{(-\lambda t^\alpha)^j}{\Gamma(\alpha(j+k)+1)} = \frac{(\lambda t^\alpha)^k}{k!} E_\alpha^{(k)}(-\lambda t^\alpha), \\ k &= 0, 1, 2, \dots, t \geq 0, \quad 0 < \alpha < 1 \end{aligned}$$

where  $E_\alpha(z)$  is the Mittag-Leffler function evaluated at  $z = -\lambda t^\alpha$ , and  $E_\alpha^{(k)}(z)$  is the  $k$ -th derivative of  $E_\alpha(z)$  evaluated at  $z = -\lambda t^\alpha$ .

Note that (see [13] and [17])

$$(18) \quad \mathbb{E}N_\alpha(t) = \int_0^\infty \mathbb{E}N(u) f_\alpha(t, u) du = \lambda t^\alpha / \Gamma(1 + \alpha),$$

where  $f_\alpha(t, u)$  is given by (13), and using the method of the bivariate integration by parts proposed in [15], [16], Leonenko, Meerschaert and Sikorskii ([17]) showed that

$$\begin{aligned} (19) \quad \text{Cov}(N_\alpha(t), N_\alpha(s)) &= \int_0^\infty \int_0^\infty \text{Var}N(1) \min(u, v) H_{t,s}(du, dv) \\ &\quad + (\mathbb{E}N(1))^2 \text{Cov}(Y_\alpha(t), Y_\alpha(s)) \\ &= \frac{\lambda(\min(t, s))^\alpha}{\Gamma(1 + \alpha)} + \lambda^2 \text{Cov}(Y_\alpha(t), Y_\alpha(s)), \end{aligned}$$

where  $\text{Cov}(Y_\alpha(t), Y_\alpha(s))$  is given in (14), and

$$H_{t,s}(u, v) = \mathbb{P}(Y_\alpha(t) < u, Y_\alpha(s) < v).$$

In this case (see again [23]),  $N_\alpha(t), t \geq 0$ , is a renewal process with Mittag-Leffler waiting times between the events

$$\begin{aligned} (20) \quad N_\alpha(t) &= \max \{n : T_1 + \dots + T_n \leq t\} = \sum_{j=1}^\infty \mathbb{I}\{T_1 + \dots + T_j \leq t\} \\ &= \sum_{j=1}^\infty \mathbb{I}\{U_j \leq G_\alpha(t)\}, \quad t \geq 0, \end{aligned}$$

where  $\{T_j\}, j = 1, 2, \dots$  are iid random variables with the strictly monotone Mittag-Leffler distribution function

$$F_\alpha(t) = \mathbb{P}(T_j \leq t) = 1 - E_\alpha(-\lambda t^\alpha), \quad t \geq 0, \quad 0 < \alpha < 1, \quad j = 1, 2, \dots,$$

and

$$G_\alpha(t) = \mathbb{P}(T_1 + \dots + T_j \leq t) = \int_0^t h^{(j)}(x) dx.$$

Here we denote an indicator as  $I\{\cdot\}$  and  $U_j, j = 1, 2, \dots$  are an iid uniformly distributed on  $[0, 1]$  random variables.  $h^{(j)}(x)$  is the pdf of  $j$ -th convolution of the Mittag-Leffler distributions which is known as the generalized Erlang distribution and it is of the form

$$h^{(j)}(x) = \alpha \lambda^j \frac{x^{j\alpha-1}}{(j-1)!} E_\alpha^{(j)}(-\lambda x^\alpha), \quad \alpha \in (0, 1), \quad x > 0.$$

Note that  $S(t) = (t/L_\alpha(1))^\alpha$  is a distribution, and

$$P(T_j > t) = Ee^{-\lambda S(t)}, \quad t \geq 0.$$

If  $h_1(t, x)$  and  $h_2(t, x)$  are densities (if they exist) of  $Y^{(1)}(t), t \geq 0$  and  $Y^{(2)}(t), t \geq 0$ , resp., then (assuming that they are independent), we have the following:

$$P(N_{FP}(t_1, t_2) = k) = \int_0^\infty \int_0^\infty P(N(x_1, x_2) = k) h_1(t_1, x_1) h_2(t_2, x_2) dx_1 dx_2, \\ k = 0, 1, 2, \dots$$

In particular, for an inverse stable subordinator, we obtain

$$(21) \quad P(N_{\alpha_1, \alpha_2}(t_1, t_2) = k) = \int_0^\infty \int_0^\infty \frac{e^{-\lambda x_1 x_2} (\lambda x_1 x_2)^k}{k!} f_{\alpha_1}(t_1, x_1) f_{\alpha_2}(t_2, x_2) dx_1 dx_2, \\ k = 0, 1, 2, \dots,$$

where  $f_\alpha(t, x)$  is given by (13). In other words,

$$(22) \quad P(N_{\alpha_1, \alpha_2}(t_1, t_2) = k) \\ = \frac{t_1 t_2 \lambda^k}{\alpha_1 \alpha_2 k!} \int_0^\infty \int_0^\infty e^{-\lambda x_1 x_2} x_1^{k-1-\frac{1}{\alpha_1}} x_2^{k-1-\frac{1}{\alpha_2}} g_{\alpha_1}(t_1 x_1^{-\frac{1}{\alpha_1}}) g_{\alpha_2}(t_2 x_2^{-\frac{1}{\alpha_2}}) dx_1 dx_2, \\ = \frac{\lambda^k}{k! t_1 t_2} \int_0^\infty \int_0^\infty e^{-\lambda x_1 x_2} x_1^{k+\frac{1}{\alpha_1}} x_2^{k+\frac{1}{\alpha_2}} W_{-\alpha_1, 0}\left(-\frac{x_1}{t_1^{\alpha_1}}\right) W_{-\alpha_2, 0}\left(-\frac{x_2}{t_2^{\alpha_2}}\right) dx_1 dx_2 \\ (t_1, t_2) \in \mathbb{R}_+^2, \quad k = 0, 1, 2, \dots,$$

where the Wright's generalized Bessel function is defined by (12), and

$$(23) \quad g_{\alpha_i}(t_i x_i^{-\frac{1}{\alpha_i}}) = \frac{1}{\pi} \sum_{j_i=1}^{\infty} (-1)^{j_i+1} \frac{\Gamma(\alpha_i j_i + 1)}{j_i!} \frac{x_i^{j_i+1/\alpha_i}}{t_i^{\alpha_i j_i+1}} \sin(\pi j_i \alpha_i), \quad i = 1, 2.$$

Using the formula

$$\sin(\pi x) = \frac{\pi}{\Gamma(x)\Gamma(1-x)},$$

one can rewrite (23) in the form

$$g_{\alpha_i}(t_i x_i^{-\frac{1}{\alpha_i}}) = \sum_{j_i=1}^{\infty} (-1)^{j_i+1} \frac{\Gamma(\alpha_i j_i + 1)}{j_i!} \frac{x_i^{j_i+1/\alpha_i}}{t_i^{\alpha_i j_i+1}} \frac{1}{\Gamma(\pi j_i \alpha_i) \Gamma(1 - \pi j_i \alpha_i)}, \quad i = 1, 2,$$

of the Mittag-Leffler type function or the H-function of Fox, see, e.g., [9].



It is known (see, e.g., [1, Theorems 1.1. and 1.2]) that a single point random field  $M$  can be regarded as a random set of points and could be characterized either by all systems of finite-dimensional distributions as

$$P(M(B_1) = k_1, \dots, M(B_m) = k_m)$$

for all integers  $m > 0$  and all compacts  $B_1, \dots, B_m$ , or the so-called capacity functional

$$T_M(K) = 1 - P(M(K) = 0),$$

for any compact  $K$ . By results of ([1, Theorems 1.1. and 1.2]), the distribution of the field  $N_{\alpha_1, \alpha_2}$  is completely characterized by its capacity functional

(24)

$$T_{N_{\alpha_1, \alpha_2}}([0, t_1] \times [0, t_2]) = 1 - \frac{t_1 t_2}{\alpha_1 \alpha_2} \int_0^\infty \int_0^\infty e^{-\lambda x_1 x_2} \frac{g_{\alpha_1}(t_1 x_1^{-\frac{1}{\alpha_1}}) g_{\alpha_2}(t_2 x_2^{-\frac{1}{\alpha_2}})}{x_1^{1+\frac{1}{\alpha_1}} x_2^{1+\frac{1}{\alpha_2}}} dx_1 dx_2.$$

Using (18) and (19) (see [13],[15], [16],[17] for details), we obtain

$$\begin{aligned} \text{EN}_{\alpha_1, \alpha_2}(t_1, t_2) &= \int_0^\infty \int_0^\infty \text{EN}(u_1, u_2) f_{\alpha_1}(t_1, u_1) f_{\alpha_2}(t_2, u_2) du_1 du_2 \\ (25) \qquad \qquad \qquad &= \lambda t_1^{\alpha_1} t_2^{\alpha_2} / [\Gamma(1 + \alpha_1) \Gamma(1 + \alpha_2)] \end{aligned}$$

and

$$\begin{aligned} &\text{Cov}(N_{\alpha_1, \alpha_2}(t_1, t_2), N_{\alpha_1, \alpha_2}(s_1, s_2)) \\ &= \int_0^\infty \int_0^\infty \int_0^\infty \int_0^\infty \text{Cov}\{N(u_1, v_1) N(u_2, v_2)\} \cdot H_{t_1, s_1}(du_1, dv_1) H_{t_2, s_2}(du_2, dv_2) \\ &\quad + \int_0^\infty \int_0^\infty \int_0^\infty \int_0^\infty \text{EN}(u_1, v_1) N(u_2, v_2) \cdot H_{t_1, s_1}(du_1, dv_1) H_{t_2, s_2}(du_2, dv_2) \\ &\quad - \lambda^2 \text{EY}_{\alpha_1}^{(1)}(t_1) \text{EY}_{\alpha_1}^{(1)}(s_1) \text{EY}_{\alpha_2}^{(2)}(t_2) \text{EY}_{\alpha_2}^{(2)}(s_2). \end{aligned}$$

We have

(26)

$$\begin{aligned} &\text{Cov}(N_{\alpha_1, \alpha_2}(t_1, t_2), N_{\alpha_1, \alpha_2}(s_1, s_2)) \\ &= \int_0^\infty \int_0^\infty \int_0^\infty \int_0^\infty \text{Var}N(1, 1) \{\min(u_1, v_1) \min(u_2, v_2)\} \cdot H_{t_1, s_1}(du_1, dv_1) H_{t_2, s_2}(du_2, dv_2) \\ &\quad + \lambda^2 (\text{EY}_{\alpha_1}^{(1)}(t_1) \text{Y}_{\alpha_1}^{(1)}(s_1) \text{EY}_{\alpha_2}^{(2)}(t_2) \text{Y}_{\alpha_2}^{(2)}(s_2) - \text{EY}_{\alpha_1}^{(1)}(t_1) \text{EY}_{\alpha_1}^{(1)}(s_1) \text{EY}_{\alpha_2}^{(2)}(t_2) \text{EY}_{\alpha_2}^{(2)}(s_2)) \\ &= \lambda (\min(t_1, s_1))^{\alpha_1} (\min(t_2, s_2))^{\alpha_2} / [\Gamma(1 + \alpha_1) \Gamma(1 + \alpha_2)] \\ &\quad + \lambda^2 \left[ \frac{1}{\Gamma(1 + \alpha_1) \Gamma(\alpha_1) \Gamma(1 + \alpha_2) \Gamma(\alpha_2)} \int_0^{\min(t_1, s_1)} \int_0^{\min(t_2, s_2)} ((t_1 - \tau_1)^{\alpha_1} + (s_1 - \tau_1)^{\alpha_1}) \right. \\ &\quad \left. \times ((t_2 - \tau_2)^{\alpha_2} + (s_2 - \tau_2)^{\alpha_2}) \tau_1^{\alpha_1 - 1} \tau_2^{\alpha_2 - 1} d\tau_1 d\tau_2 - \frac{(s_1 t_1)^{\alpha_1} (s_2 t_2)^{\alpha_2}}{\Gamma^2(1 + \alpha_1) \Gamma^2(1 + \alpha_2)} \right], \end{aligned}$$

$$(t_1, t_2), (s_1, s_2) \in \mathbb{R}_+^2,$$

in particular

(27)

$$\text{Var}N_{\alpha_1, \alpha_2}(t_1, t_2) = \lambda t_1^{\alpha_1} t_2^{\alpha_2} C_1(\alpha_1, \alpha_2) + \lambda^2 t_1^{2\alpha_1} t_2^{2\alpha_2} C_2(\alpha_1, \alpha_2), \quad (t_1, t_2) \in \mathbb{R}_+^2,$$

where

$$C_1(\alpha_1, \alpha_2) = \frac{1}{\Gamma(1 + \alpha_1)\Gamma(1 + \alpha_2)}; \quad C_2(\alpha_1, \alpha_2) = \frac{1}{\alpha_1 \alpha_2 \Gamma(2\alpha_1)\Gamma(2\alpha_2)} - \frac{1}{(\alpha_1 \alpha_2)^2 \Gamma^2(\alpha_1)\Gamma^2(\alpha_2)}$$

We can summarize our results in the following

**Theorem.** *Let  $N_{\alpha_1, \alpha_2}(t_1, t_2), (t_1, t_2) \in \mathbb{R}_+^2$ , be a fractional Poisson process defined by (17). Then*

- i)  $P(N_{\alpha_1, \alpha_2}(t_1, t_2) = k), k = 0, 1, 2, \dots$  is given by (22);
- ii)  $EN_{\alpha_1, \alpha_2}(t_1, t_2), \text{Var}N_{\alpha_1, \alpha_2}(t_1, t_2)$  and  $\text{Cov}(N_{\alpha_1, \alpha_2}(t_1, t_2), N_{\alpha_1, \alpha_2}(s_1, s_2))$  are given by (25), (27), (26) respectively;
- iii) the following long-range dependence property holds:

$$(28) \quad \frac{N_{\alpha_1, \alpha_2}(t_1, t_2)}{t_1^{2\alpha_1} t_2^{2\alpha_2}} - \frac{\lambda}{t_1^{\alpha_1} t_2^{\alpha_2} [\Gamma(1 + \alpha_1)\Gamma(1 + \alpha_2)]} \rightarrow 0, \quad \alpha_1 \in (0, 1), \quad \alpha_2 \in (0, 1),$$

in probability as  $t_1, t_2 \rightarrow \infty$ .

*Remark.* For  $d = 1$  the definition of Hurst index for renewal processes is discussed in [7]. Following the ideas of this paper the Hurst index of the fractional Poisson random field in  $d = 2$  can be defined as follows:

$$H = \inf \left\{ \beta : \limsup_{T \rightarrow \infty} \frac{\text{Var}N_{\alpha_1, \alpha_2}(T, T)}{T^\beta} < \infty \right\},$$

thus  $H = 2(\alpha_1 + \alpha_2)$  and lies in the interval  $0 < H < 4$ .

## 5. ALTERNATIVE FORMS OF THE FRACTIONAL POISSON-LIKE RANDOM FIELD

We present some ideas for different versions of fractional Poisson-like random fields.

Similar to (20) one can define a point random field

$$(29) \quad N_F^{(0)}(t_1, t_2) = \sum_{j=1}^{\infty} \mathbf{I}\{U_j \in [0, G_{\alpha_1}(t_1)] \times [0, G_{\alpha_2}(t_2)]\}, \quad (t_1, t_2) \in \mathbb{R}_+^2,$$

where  $U_j = (U_j^{(1)}, U_j^{(2)}), j = 1, 2, \dots$  is a sequence of iid random vectors on  $[0, 1]^2 \subset \mathbb{R}_+^2$  with the uniform distribution, that is, for any Borel  $A \subset [0, 1]^2$

$$P(U_j \in A) = |A|,$$

and

$$0 < \alpha_i < 1, \quad i = 1, 2.$$

This definition is motivated by Ivanoff, Merzbach [12], Herbin, Merzbach [10] and Merzbach, Shaki [25]. Formula (29) is useful for computations as well as

simulation. Notice that all components  $U_j^{(1)}, U_j^{(2)}, j = 1, 2, \dots$  are independent, and

$$\begin{aligned} \mathbb{I}\{U_j \in [0, G_{\alpha_1}(t_1)] \times [0, G_{\alpha_2}(t_2)]\} &= \mathbb{I}\{G_{\alpha_1}(U_j^{(1)}) \leq t_1, G_{\alpha_1}(U_j^{(2)}) \leq t_2\} \\ &= \mathbb{I}\{G_{\alpha_1}(U_j^{(1)}) \leq t_1\} \mathbb{I}\{G_{\alpha_1}(U_j^{(2)}) \leq t_1\} \\ &= \mathbb{I}\{T_1^{(1)} + \dots + T_j^{(1)} \leq t_1\} \mathbb{I}\{T_1^{(2)} + \dots + T_j^{(2)} \leq t_2\}, \end{aligned}$$

where  $\{T_j^{(i)}\}, j = 1, 2, \dots$  are iid random variables with the strictly monotone Mittag-Leffler distribution function

$$F_{\alpha_i}(t) = \mathbb{P}\left(T_j^{(i)} \leq t\right) = 1 - E_{\alpha_i}(-\lambda t^{\alpha_i}), \quad t \geq 0, \quad 0 < \alpha_i < 1, \quad j = 1, 2, \dots, \quad i = 1, 2.$$

Thus, the fractional Poisson-like random field (29) can be written as

$$\tilde{N}_F(t_1, t_2) = \sum_{j=1}^{\infty} \mathbb{I}\{T_1^{(1)} + \dots + T_j^{(1)} \leq t_1\} \mathbb{I}\{T_1^{(2)} + \dots + T_j^{(2)} \leq t_2\}, \quad (t_1, t_2) \in \mathbb{R}_+^2,$$

and  $\mathbb{P}\left\{N_F^{(0)}(t_1, t_2) = k\right\}$  can be computed, while its capacity functional  $T_{\tilde{N}}(K) = 1 - \mathbb{P}(\tilde{N}_F(K) = 0)$  is different from (24).

One can also define the Poisson-like random fields  $N_F^{(k)}(A), k \in \{1, 2, 3\}$  by their capacity functionals

$$T_{N^{(1)}}(A) = 1 - \mathbb{P}(N_F^{(1)}(A) = 0) = 1 - \frac{1}{E_{\alpha}(\lambda |A|)}, \quad 0 < \alpha \leq 1, \quad \lambda > 0,$$

or

$$T_{N^{(2)}}(A) = 1 - \mathbb{P}(N_F^{(2)}(A) = 0) = 1 - E_{\alpha}(-\lambda |A|^{\alpha}),$$

for any Borel set  $A \subset \mathbb{R}^2$  with finite Lebesgue measure  $|A| < \infty$ . The above constructions coincide with the definition of the homogeneous Poisson random field for  $\alpha = 1$ , but for  $0 < \alpha < 1$  many basic properties (such as independence of increments) are lost and we even do not know if there exist Poisson-like fields with the properties

$$\mathbb{P}(N_F^{(1)}(A) = k) = \frac{(\lambda |A|)^k}{E_{\alpha}(\lambda |A|)\Gamma(\alpha k + 1)}, \quad k = 0, 1, 2, \dots$$

or

$$\mathbb{P}(N_F^{(2)}(A) = k) = \frac{(\lambda |A|^{\alpha})^k}{k!} \sum_{r=0}^{\infty} \frac{(r+k)!}{r!} \frac{(\lambda |A|^{\alpha})^r}{\Gamma(\alpha(r+k) + 1)}, \quad k = 0, 1, 2, \dots,$$

or

(30)

$$\mathbb{P}(N_F^{(3)}(A) = k) = \frac{(-1)^k}{k!} \sum_{r=0}^{\infty} \frac{(-\lambda |A|^{\alpha})^r}{r!} \frac{(r+k)!}{r!} \frac{\Gamma(\alpha r + 1)}{\Gamma(\alpha r + 1 - k)}, \quad k = 0, 1, 2, \dots,$$

see however [3], [4],[26] for some discussion of the problem for  $d = 1$ .

These problems will be discussed elsewhere.

## ACKNOWLEDGMENTS

The authors wish to thank both referees for valuable comments and corrections which improve the presentation and, in particular, one of them who proposed formula (30).

## REFERENCES

- [1] A. Baddeley, Spatial point processes and their applications. Stochastic geometry, 1–75, Lecture Notes in Math., 1892, Springer, Berlin, 2007
- [2] E. Barkai, Fractional Fokker-Planck equation, solution, and application, Physical Review E 63 (2001), 046118.
- [3] L. Beghin and E. Orsingher, Fractional Poisson processes and related planar random motions, Electron. J. Probab. 14 (2009), no. 61, 1790–1827.
- [4] L. Beghin and E. Orsingher, Poisson-type processes governed by fractional and higher-order recursive differential equations. Electron. J. Probab. 15 (2010), no. 22, 684–709.
- [5] N. H. Bingham, Limit theorems for occupation times of Markov processes, Z. Wahrscheinlichkeitstheorie und Verw. Gebiete 17 (1971), 1–22.
- [6] D. O. Cahoy, V. V. Uchaikin, and W. A. Woyczynski, Parameter estimation for fractional Poisson processes, J. Statist. Plann. Inference 140 (2010), no. 11, 3106–3120.
- [7] D.J.Daley, The Hurst index for a long-range dependent renewal processes, Annals of Probability, 27 (1999), no 4, 2035–2041.
- [8] M.M. Djrbashian, Harmonic Analysis and Boundary Value Problems in the Complex Domain, Birkhauser Verlag, Basel, 1993.
- [9] H. J. Haubold, A. M. Mathai, and R. K. Saxena, Mittag-Leffler functions and their applications. J. Appl. Math. 2011, Art. ID 298628, 51 pp.
- [10] E. Herbin and E. Merzbach, A set-indexed fractional Brownian motion, J. Theor. Probab. 19(2) (2006), 337–364.
- [11] G. Ivanoff and E. Merzbach, Set-indexed Martingales, Chapman & Hall, 2000.
- [12] B.G. Ivanoff and E. Merzbach, What is a multi-parameter renewal process?, Stochastics 78 (2006), no. 6, 411–441.
- [13] J.Janczura and A. Wylomanska, Subdynamics of financial data from fractional Fokker-Planck equation, Acta Physica Polonica B 40 (2009), 1341–1351.
- [14] N. Laskin, Fractional Poisson process. Chaotic transport and complexity in classical and quantum dynamics, Commun. Nonlinear Sci. Numer. Simul. 8 (2003), no. 3-4, 201–213.
- [15] N.N. Leonenko, M.M. Meerschaert and A. Sikorskii, Fractional Pearson diffusions, Journal of Mathematical Analysis and Applications, 403 (2013), 532–246.
- [16] N.N. Leonenko, M.M. Meerschaert and A. Sikorskii, Correlation structure of fractional Pearson diffusions, Computers and Mathematics and Applications, in press, doi:10.1016/j.camwa.2013.01.009
- [17] N.N. Leonenko, M.M. Meerschaert and A. Sikorskii, Correlation structure of time changed Levy processes, preprint (2013)
- [18] N.N. Leonenko, M.D. Ruiz-Medina and M.S. Taqqu, Fractional elliptic, hyperbolic and parabolic random fields, Electronic Journal of Probability 16 (2011), 1134–1172.
- [19] F. Mainardi, R. Gorenflo, and E. Scalas, A fractional generalization of the Poisson processes, Vietnam J. Math. 32, 2004, Special Issue, 53–64.
- [20] F. Mainardi, R. Gorenflo, and A. Vivoli, Renewal processes of Mittag-Leffler and Wright type, Fract. Calc. Appl. Anal. 8 (2005), no. 1, 7–38.
- [21] F. Mainardi, R. Gorenflo, and A. Vivoli, Beyond the Poisson renewal process: a tutorial survey. J. Comput. Appl. Math. 205 (2007), no. 2, 725–735.

- [22] M. M. Meerschaert and A. Sikorskii, *Stochastic Models for Fractional Calculus*, De Gruyter, Berlin/Boston, 2012.
- [23] M. M. Meerschaert, E. Nane, and P. Vellaisamy, The fractional Poisson process and the inverse stable subordinator, *Electron. J. Probab.* 16 (2011), no. 59, 1600-1620.
- [24] E. Merzbach and D. Nualart, A characterization of the spatial Poisson process and changing time, *Ann.Probab.* 14 (1986), no. 4, 1380-1390.
- [25] E. Merzbach and Y.Y. Shaki, Characterizations of multiparameter Cox and Poisson processes by the renewal property, *Stat. Probab. Letters* 78 (2008), 637-642.
- [26] E. Orsingher and Polito, The space-fractional Poisson process, *Stat. Probab. Letters* 82 (2012), 852-858.
- [27] I. Podlubny, *Fractional differential equations. An introduction to fractional derivatives, fractional differential equations, to methods of their solution and some of their applications*, Mathematics in Science and Engineering, 198, Academic Press, Inc., San Diego, CA, 1999.
- [28] O.N. Repin and A.I. Saichev, Fractional Poisson law, *Radiophysics and Quantum Electronics* 43 (2000), no. 9, 738-741.
- [29] D. Stoyan, W.S. Kendall, and J. Mecke, *Stochastic Geometry and its Applications*, Wiley, New York, 1995.
- [30] M. Veillette, M. S. Taqqu, and S. Murad, Numerical computation of first passage times of increasing Lévy processes, *Methodol. Comput. Appl. Probab.* 12 (2010), no. 4, 695–729.
- [31] M. Veillette and M. S. Taqqu, Using differential equations to obtain joint moments of first-passage times of increasing Lévy processes, *Statist. Probab. Lett.* 80 (2010), no. 7-8, 697–705.
- [32] X.-T. Wang and Z.-X. Wen, Poisson fractional processes, *Chaos Solitons Fractals* 18 (2003), no. 1, 169–177.
- [33] X.-T. Wang, Z.-X. Wen, and S.-Y. Zhang, Fractional Poisson process. II, *Chaos Solitons Fractals* 28 (2006), no. 1, 143–147.
- [34] X.-T. Wang, S.-Y. Zhang, and S. Fan, Nonhomogeneous fractional Poisson processes. *Chaos Solitons Fractals* 31 (2007), no. 1, 236–241.

SCHOOL OF MATHEMATICS, CARDIFF UNIVERSITY, SENGHENNYDD ROAD, CARDIFF CF24 4AG, UK

*E-mail address:* LeonenkoN@Cardiff.ac.uk

DEPARTMENT OF MATHEMATICS, BAR-ILAN UNIVERSITY, 52900 RAMAT-GAN, ISRAEL

*E-mail address:* merzbach@macs.biu.ac.il